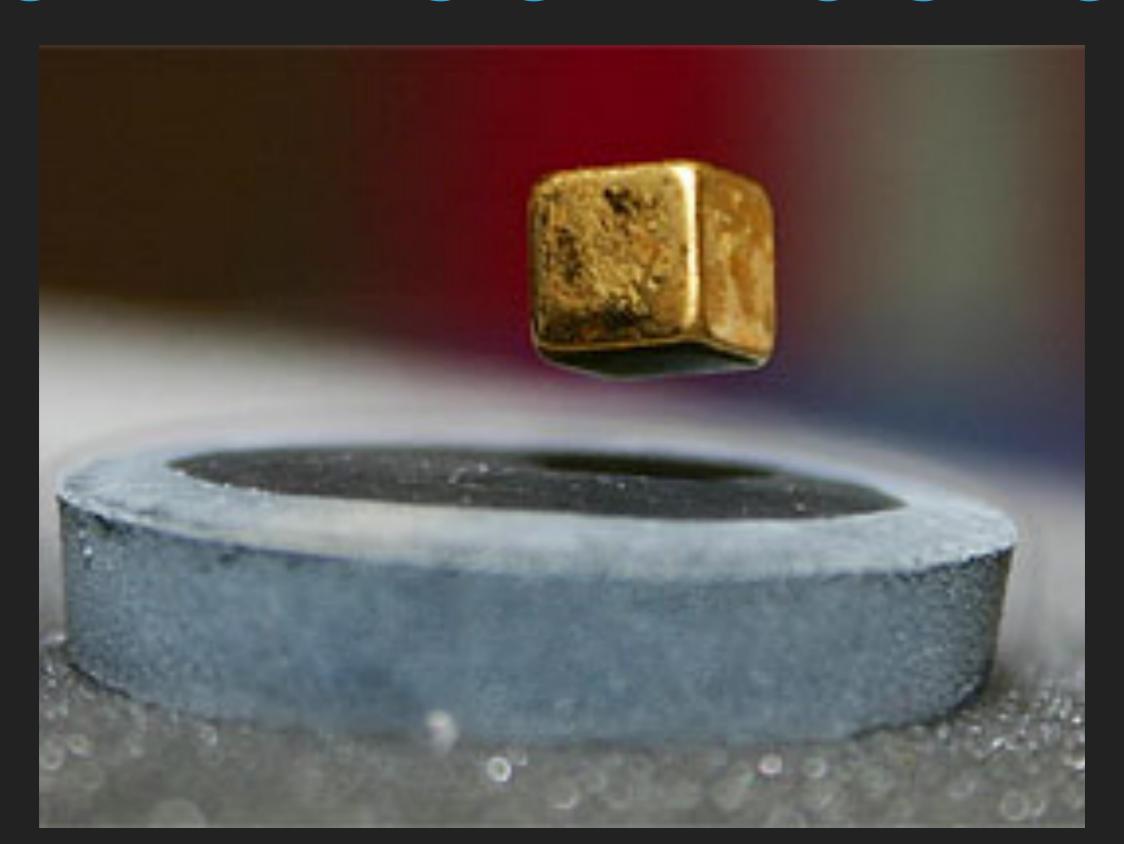
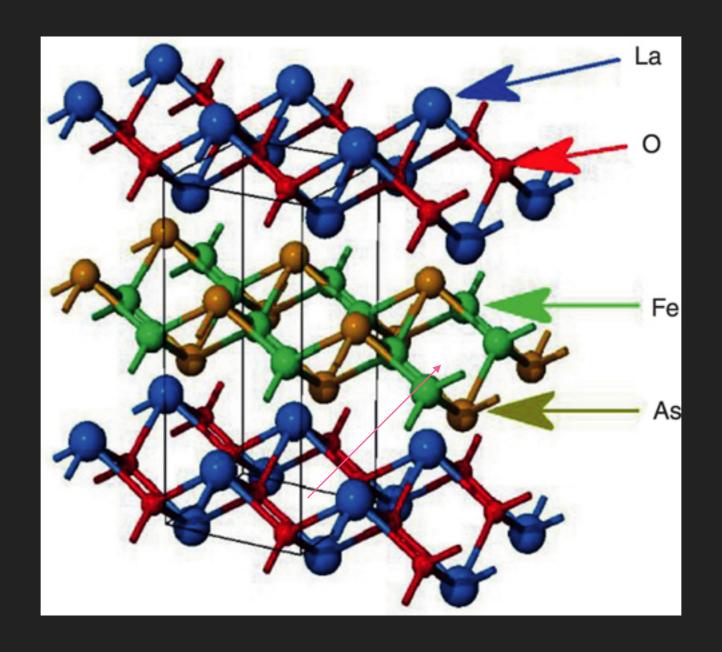
#### ROBERT HELLING (LMU)

# HOLOGRAPHIC THEORIES: AN INVITATION

# SUPER CONDUCTOR



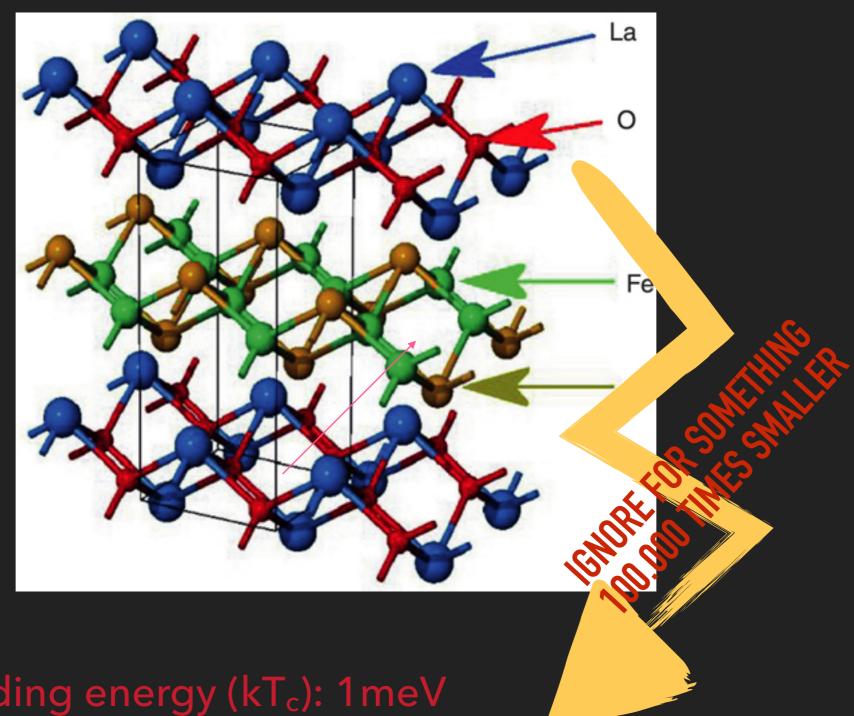
## MORE REALISTICALLY...



Cooper pair binding energy (kT<sub>c</sub>): 1meV

Atomic binding energy (Coulomb): 100eV

### MORE REALISTICALLY...



Cooper pair binding energy (kT<sub>c</sub>): 1meV

Atomic binding energy (Coulomb): 100eV



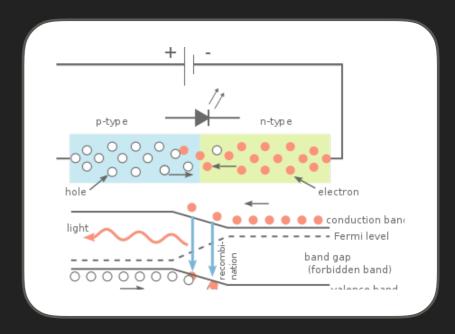
#### **BCS IS AN EFFECTIVE THEORY**

Microscopic (UV) theory is know:

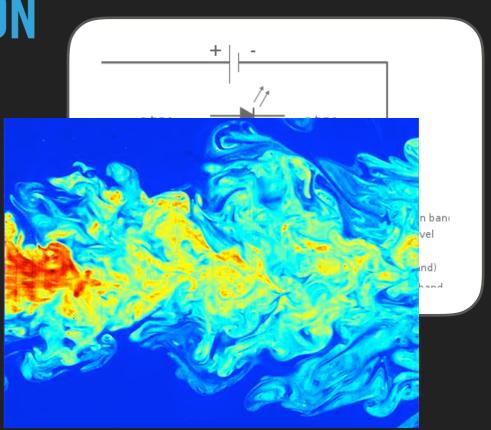
$$H = -\sum_{i} \Delta_{i} - \sum_{i,a} \frac{Z_{a}}{|R_{a} - r_{i}|} + \sum_{i < j} \frac{1}{|r_{i} - r_{j}|}$$

- Superconductivity is a phenomenon of the IR end of the renormalization group flow
- There is little control over the RG flow, so BCS as an "effective theory"

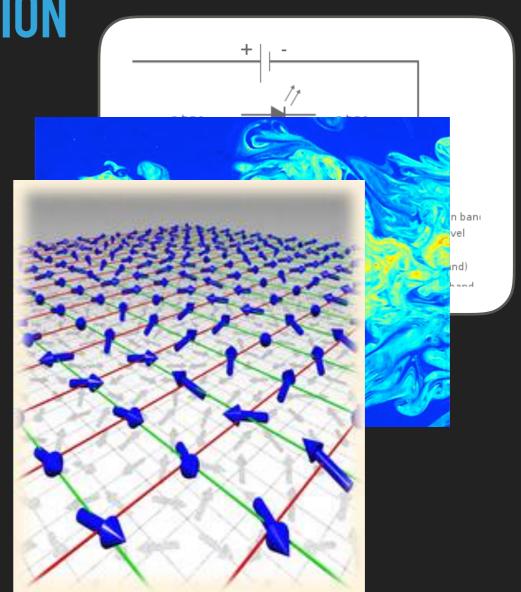
Fermi Liquid



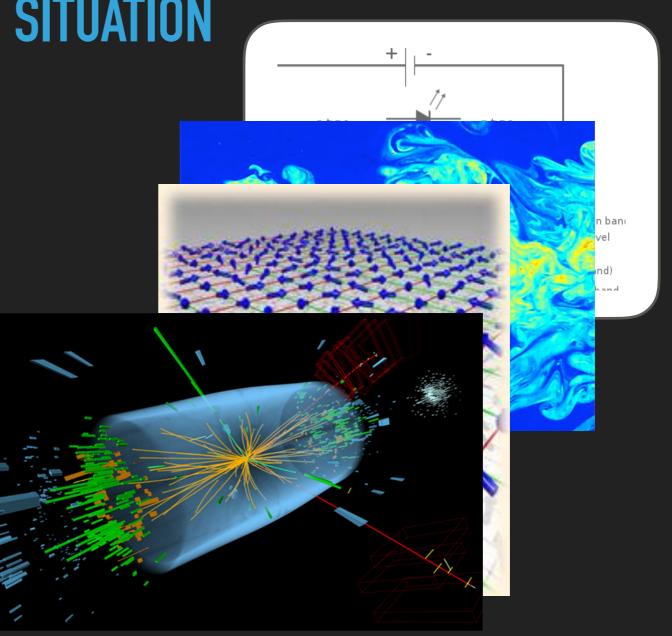
- Fermi Liquid
- Hydrodynamics



- Fermi Liquid
- Hydrodynamics
- Ferromagnets



- Fermi Liquid
- Hydrodynamics
- Ferromagnets
- Standard Model



#### WHAT THE PHYSICIST NEEDS...

- What is thought for is a "simpler" theory that makes qualitative and quantitative predictions for low energy phenomena (e.g. transport properties like conductivity)
- In the end, in comes down to guessing a quantum (field) theory (a model).
- To have a starting point, we need a language to write down the QFT.
- Output are n-point (Wightman) functions, e.g.

$$\sigma_{\alpha\beta}(\omega) = \frac{ine^2}{m\omega} \delta_{\alpha\beta} + \frac{1}{\hbar\omega V} \int_0^\infty dt \, e^{i\omega t} \langle [j_{\alpha}(t), j)_{\beta}(t)] \rangle$$

Free fields

- Free fields
  - plus perturbation expansion

- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles

- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles
  - weak coupling

- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles
  - weak coupling
- Strong symmetry constraints

- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles
  - weak coupling
- Strong symmetry constraints
  - ▶ 2D CFT (e.g. Kondo effect)

- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles
  - weak coupling
- Strong symmetry constraints
  - 2D CFT (e.g. Kondo effect)
  - minimal models

- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles
  - weak coupling
- Strong symmetry constraints
  - 2D CFT (e.g. Kondo effect)
  - minimal models
- Lattice methods

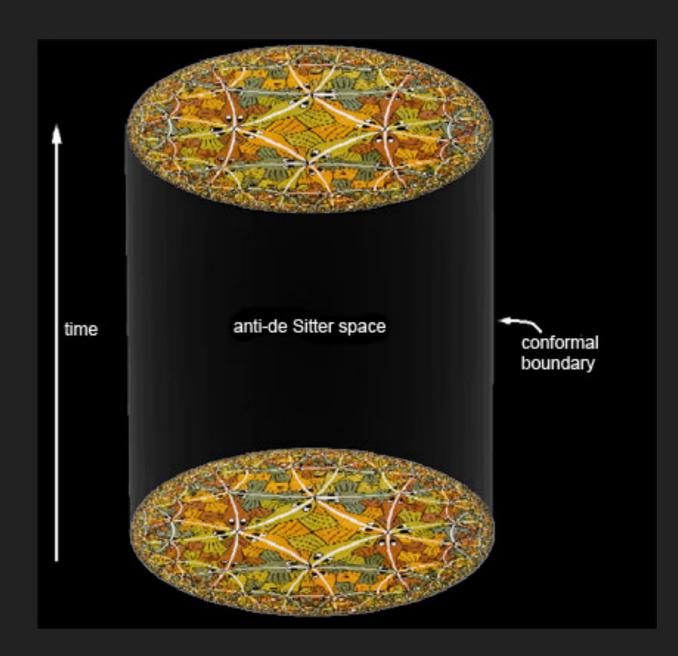
- Free fields
  - plus perturbation expansion
  - Lagrangian, quasi-particles
  - weak coupling
- Strong symmetry constraints
  - 2D CFT (e.g. Kondo effect)
  - minimal models
- Lattice methods
- ▶ Holography

#### ANTI-DE-SITTER-SPACE

$$ds^2 = \frac{L^2}{r^2} \left( dr^2 + \eta_{\mu\nu} dx^{\mu} dx^n \right)$$

Classical bulk action
with asymptotic prescription
computes
n-point functions of
QFT on the boundary

$$\left\langle e^{\mathcal{O}(\phi_0)} \right\rangle_{QFT} = e^{i\mathcal{S}_{bulk}} | "_{\phi \to \phi_0} "$$



#### **ADS/CFT DICTIONARY**

- Pick a classical bulk theory (gravity + other fields (scalars, vectors, etc.). Require solutions with asymptotic  $(r \rightarrow 0)$  AdS geometry (e.g. AdS-Black hole for T>0)
- For each field, there is a corresponding operator in the boundary theory, e.g.

$$A_{\mu}(x,r) \leftrightarrow J_{m}(x)$$
  
 $g_{\mu\nu}(x,r) \leftrightarrow T_{mn}(x)$ 

#### AN EXAMPLE: A SCALAR FIELD

For  $\phi(x,r)=\phi(r)e^{ikx}$  Klein-Gordon equation becomes

$$-r^{d+1}\partial_r (r^{-d+1}\partial_r \phi) + (k^2r^2 + m^2L^2) \phi = 0$$

Near the r=0 boundary,

$$\phi(r) = \left(\frac{r}{L}\right)^{\Delta_{-}} \left(\phi_{0}(x) + \cdots\right) + \left(\frac{r}{L}\right)^{\Delta_{+}} \left(\phi_{1}(x) + \cdots\right)$$

$$\Delta_{\pm} = \frac{d}{2} \pm \sqrt{\frac{d^2}{4} + m^2 L^2}$$

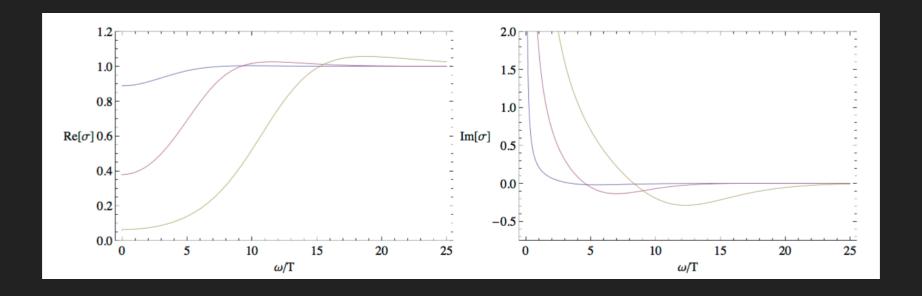
#### RECIPE

$$\phi(r) = \left(\frac{r}{L}\right)^{\Delta_{-}} \left(\phi_{0}(x) + \cdots\right) + \left(\frac{r}{L}\right)^{\Delta_{+}} \left(\phi_{1}(x) + \cdots\right)$$

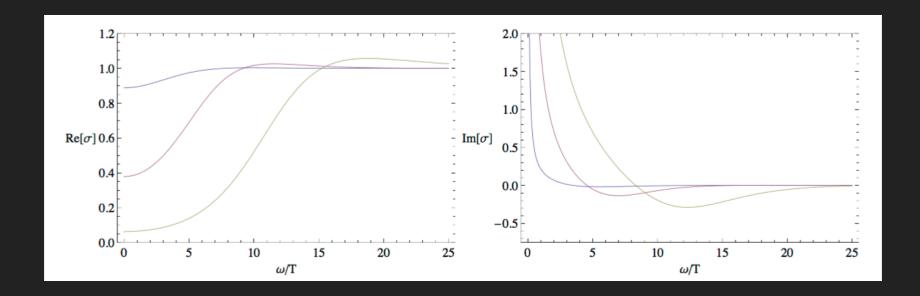
- ullet Fix  $\phi_0(x)$  using the boundary source (test function)
- Impose a boundary condition at large r (typically: purely in-going at BH horizon)
- Solve the bulk equations of motion
- Read off  $\phi_1(x)$  as linear response

A similar calculation for a gauge field (dual to current):

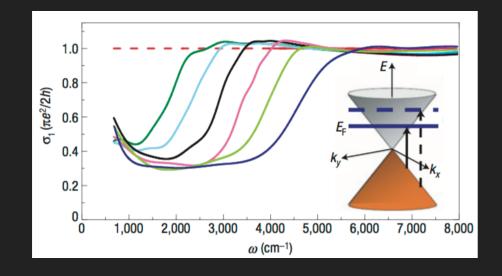
A similar calculation for a gauge field (dual to current):

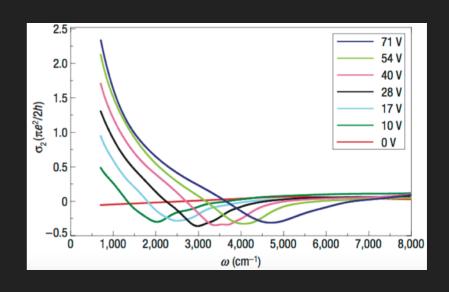


A similar calculation for a gauge field (dual to current):

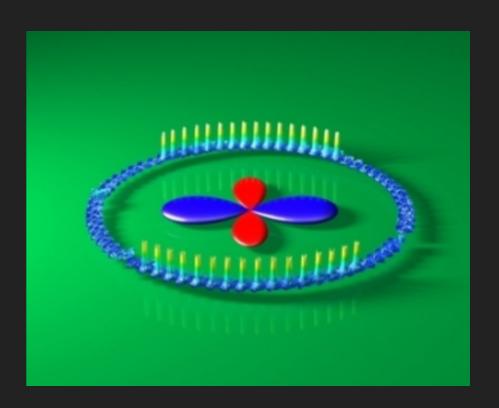


Compare to graphene (measured):

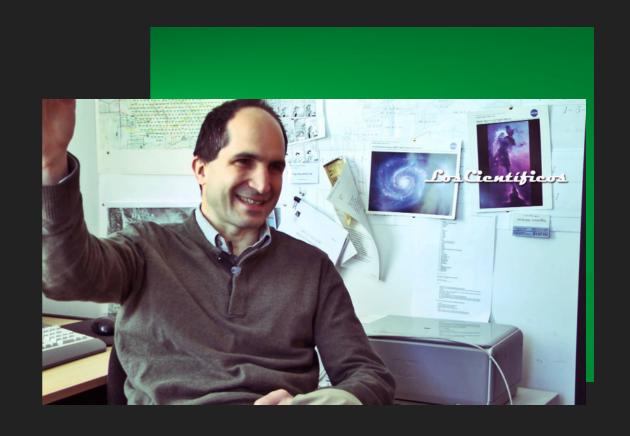




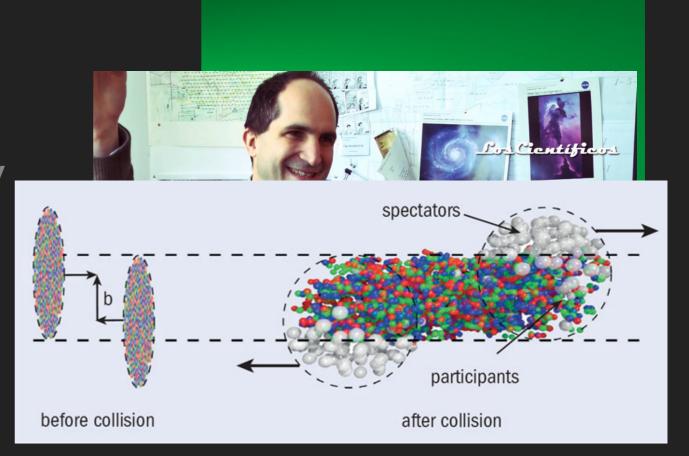
d-wave superconductors



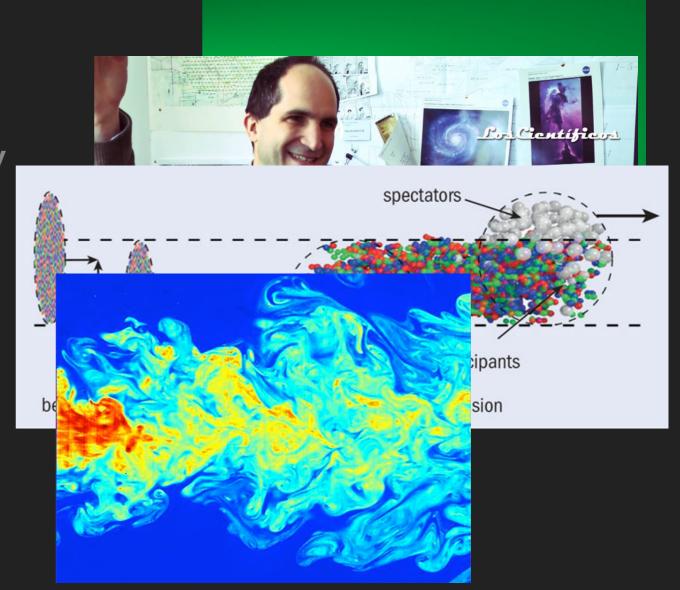
- d-wave superconductors
- $ightharpoonup \mathcal{N} = 4 \; SU(N) \; \mathrm{YM} \; \mathrm{theory}$



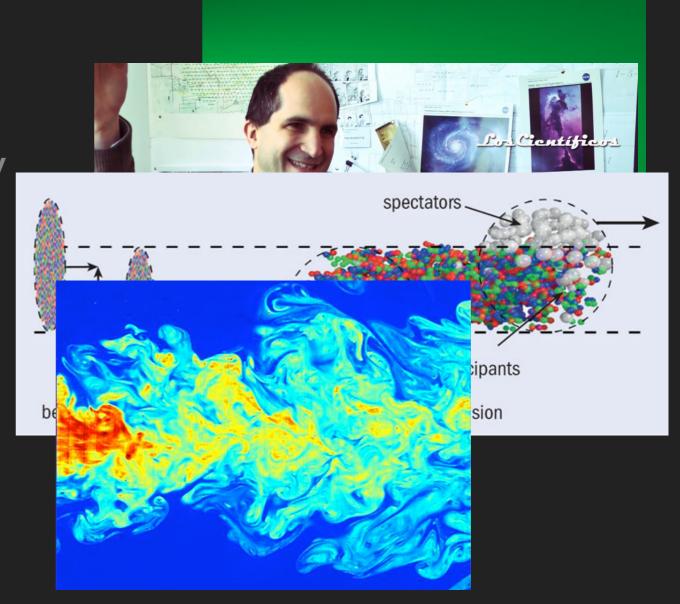
- d-wave superconductors
- $ightharpoonup \mathcal{N} = 4 \ SU(N) \ \mathrm{YM} \ \mathrm{theory}$
- heavy ion collisions



- d-wave superconductors
- $ightharpoonup \mathcal{N} = 4 \; SU(N) \; ext{YM theory}$
- heavy ion collisions
- hydrodynamics



- d-wave superconductors
- $\mathcal{N}=4$  SU(N) YM theory
- heavy ion collisions
- hydrodynamics
- **...**



#### MY QUESTION TO YOU

- Does this approach to correlation functions really encode a QFT (à la Wightman, Osterwalder-Schrader etc)?
- If yes, alternative route to free fields and 2d CFT
- If no, interesting to the holographic community

#### **FURTHER READING**

- Sean A. Harntoll "Lectures on holographic methods for condensed matter physics" arxiv:0903.3264 (source for conductivity plots)
- David Tong "Lectures on Holographic Conductivity" http:// www.damtp.cam.ac.uk/user/tong/ talks/zakopane.pdf
- Kostas Skenderis, Balt C. van Rees "Real-time gauge-gravity duality: Prescription, Renormalization and Examples" arXiv:0812.2909